

# Ionization of the E Region of the Ionosphere by Precipitated Electrons

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It is shown that there is a good correlation between the fluxes of electrons with energies greater than 40 keV observed by Alouette I near the L shell 2,4 and the critical frequency of the E-layer observed at Argentine Islands, which lies on the same L shell. The increase of  $f_o E$  observed agrees reasonably well with that calculated on the basis of a precipitated electron flux with characteristic energy 50 keV, of which the portion above 40 keV is observed by Alouette. It is shown that the rate of ionization by the precipitated electrons can exceed that due to solar ultraviolet radiation at times.

## Introduction

In the course of a study of the relationship between ionospheric behaviour and the precipitated electron fluxes observed by the satellite Alouette I, Gledhill *et al.* (1967) showed that there was a very significant correlation between the flux of electrons with energies greater than 40 keV, observed in the vicinity of an ionospheric station, and the critical frequency of the F2 layer there. Since electrons with such high energies have little direct effect on the atmosphere at F-region heights, depositing their energy around the 85 km level far below, it was clear that the correlations found by Gledhill *et al.* (1967) must imply that the energetic electrons were accompanied by fluxes of lower-energy particles which were the effective agents in ionizing the F-region. Further, these low-energy electrons must vary in much the same way spatially and temporally as did the fluxes above 40 keV, which were being used as indicators of their presence.

The energetic electrons actually observed by Alouette I should be effective in ionizing the D region, where their main observable influence would be to increase the absorption of radio waves and to alter the height of reflection of low frequency waves. The only routinely published ionospheric parameter which depends on conditions in the D region is  $f_{min}$ , the lowest frequency at which an echo is observed with an ionosonde. This is a notoriously unreliable parameter, which varies with transmitter efficiency, receiver gain, antenna and ground characteristics and so even with such quantities as the mains supply voltage unless this is regulated, which is often not the case at remote stations. Attempts (Haschick, 1973) to correlate  $f_{min}$  with the precipitated electron flux for Sanae ( $2^{\circ} 21' W; 70^{\circ} 18' S$ ) and for Argentine Islands ( $64^{\circ} 16' W; 65^{\circ} 15' S$ ) led to barely significant results. This may have been because no real-time telemetry was receivable while the satellite was in the immediate vicinity of Sanae; while it was later learned that the attenuator setting on the ionosonde receiver at Argentine Islands was sometimes

*Daar word aangetoon dat 'n goeie korrelasie bestaan tussen die vloed van elektrone met energie hoër as 40 keV wat deur Alouette I naby die L-skil 2,4 waargeneem is en die kritiese frekwensie van die E-laag wat by die Argentine-eilande waargeneem is en wat op dieselfde L-skil lê. Die waargenome toename van  $f_o E$  stem taamlik goed ooreen met wat bereken is op grondslag van 'n gespesifieerde elektronvloed met karakteristieke energie van 50 keV, waarvan die gedeelte bo 40 keV deur Alouette I waargeneem word. Daar word aangetoon dat die tempo van ionisasie deur die gespesifieerde elektrone soms hoër kan wees as wat ultraviolet-sonstraling veroorsaak.*

altered according to conditions, so that  $f_{min}$  loses much of its significance for correlations of this type.

Although 40 keV electrons leave only about 1/30 as much energy per unit volume at 120 km as they do at their peak near 85 km (Rees, 1963), it was felt that their flux might be more closely correlated with that of the 5 keV electrons which have their peak loss rate in the E region, than with the flux of soft electrons which would ionize the F region. An analysis was therefore made of the relationship between the critical frequency of the E layer,  $f_o E$ , observed at Argentine Islands, and the flux of electrons above 40 keV measured overhead by Alouette I.

The choice of Argentine Islands for this purpose was governed by several considerations:

- (i) the ionograms obtained there were of a uniformly good quality, and it was usually possible to scale  $f_o E$  with reasonable precision;
- (ii) the station lies within the area from which real-time telemetry was received at the South Atlantic station, so that flux measurements could be selected, which were made within an acceptable distance of the overhead point at the station;
- (iii) the station lies on the western edge of the South Atlantic Anomaly region (Vernov *et al.*, 1967), so that electron precipitation from the radiation belts would be more frequent and intense there than at other stations for which ionograms were available;
- (iv) although the station is situated near the Antarctic Peninsula ( $64^{\circ} 16' W; 65^{\circ} 15' S$ ), it is not near the auroral zone (its L value is only 2,4) and it is indeed a middle latitude station from the geomagnetic point of view (geomagnetic latitude  $-53,8^{\circ}$ ).

The analysis covers the period 1-18 November 1962. No ionograms were available for Argentine Islands from 19 to 26 November, while background counts due to Russian high-altitude nuclear explosions earlier in the year made

Table 1

Nominal local time (h)	Solar zenith angle (°)	Correlation coefficient r	Significance level	Slope (MHz <sup>4</sup> cm <sup>2</sup> s sr)	Intercept (MHz <sup>4</sup> )	$f_0 E_0$ ± (std. dev.) MHz
0200	95.00	0.59	0.07	0.216	0.254	0.71 ± 0.25
0300	91.34	0.81	0.003	0.327	2.11	1.20 ± 0.05
0400	86.57	0.43	0.2	0.362	7.35	1.65 ± 0.06
0500	81.00	0.07	0.7	0.13	18.4	2.07 ± 0.05
1700	68.95	0.91	$3 \times 10^{-5}$	1.57	40.7	2.52 ± 0.01
1800	74.91	0.85	$10^{-4}$	1.39	27.4	2.20 ± 0.03
1900	81.00	0.88	$3 \times 10^{-4}$	1.34	13.7	1.92 ± 0.02
2000	86.57	0.77	0.02	0.615	6.27	1.58 ± 0.05

the electron flux measurements rather unreliable before the end of October. During this interval, particle counts were available for morning (02-05 h LT) and evening (17-20 h) passes and there was a relatively high level of variation in the fluxes measured from pass to pass, thus giving a good chance of finding the correlation if it existed.

### Preliminary Analysis

The L value at about 200 km above Argentine Islands is 2.45. The shell  $2.26 < L < 2.63$  corresponds to a region of width about 350 km centred on this point. All telemetered passes of the satellite through this shell were considered, covering the longitude range  $275^\circ$  to  $330^\circ$ E and thus extending about 1 000 km west of Argentine Islands and 1 500 km east of it. Since trapped electrons drift along such shells, it seems reasonable to expect a good correlation over a considerably greater distance along the shell

than across it. The height of the satellite was a little above 1 000 km.

The electron counters on Alouette I have been described by McDiarmid *et al.* (1963). The count rates from the Anton 223 counter with the 40 keV threshold, the Anton 223 counter with the 250 keV threshold and the Anton 302 counter were used as described by Greener & Gledhill (1972) to derive corrected directional fluxes of electrons  $> 40$  keV and  $> 250$  keV. The latter was then subtracted from the former to give the flux between 40 and 250 keV. Several 10-second counts were recorded during each pass through the L-shell considered and the average was taken for further consideration. No account was taken of the pitch angle at which the electron flux was observed.

Copies of the ionograms from Argentine Islands were carefully examined for times within 1.5 hours of a satellite pass. The frequencies of all cusps in the vicinity of the normal  $f_0 E_0$  were very carefully scaled and plotted against

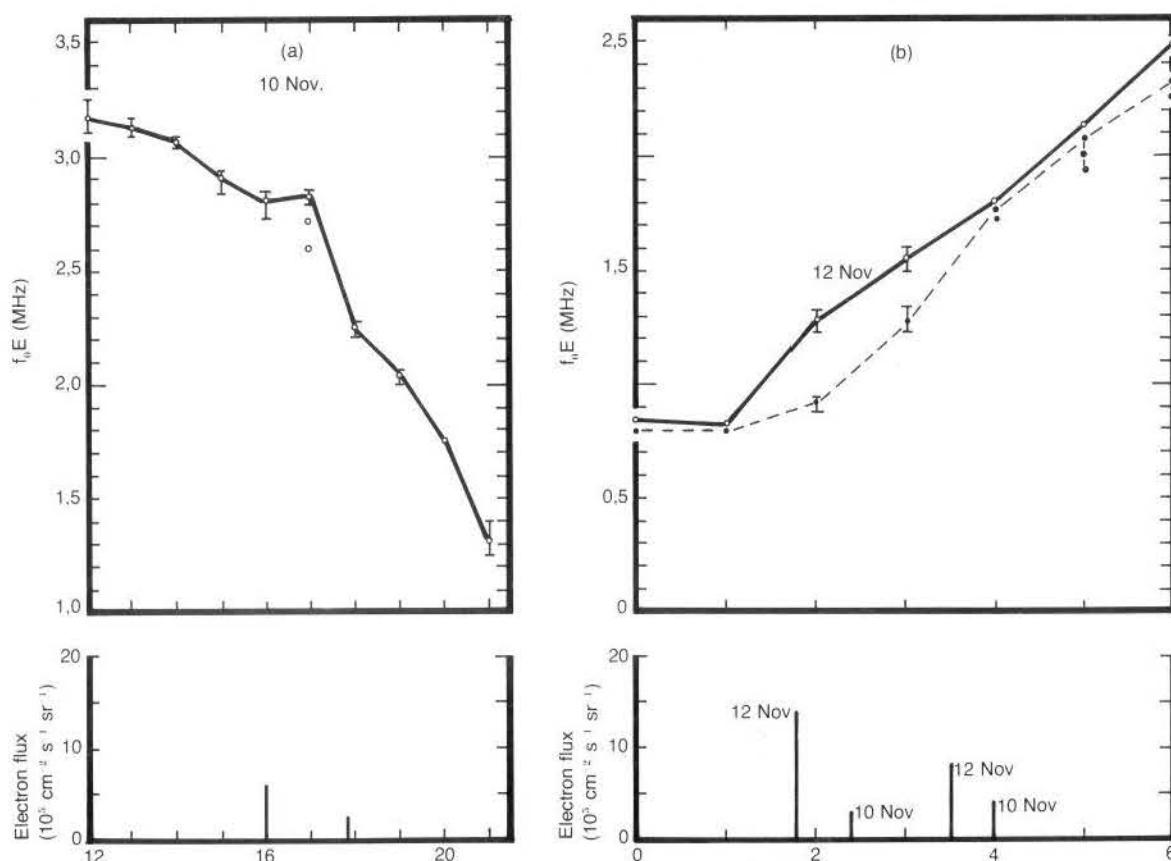


Fig. 1. Increases of  $f_0 E_0$  coinciding with high electron fluxes at Argentine Islands.

time. The limits of error within which each cusp certainly lay were recorded. The r.m.s. error limit was 0.05 MHz, though sometimes the frequency of a cusp could not be determined to better than 0.1 MHz, while sometimes it could be established to within 0.01 MHz. It soon became clear which cusp represented the ordinary E layer critical frequency. On several occasions when the electron flux measured by the satellite was high, a small digression of  $f_oE$  from its normal diurnal variation was noted. Fig. 1 shows two such occasions, *viz* when an unusually high  $f_oE$ , with two extra cusps, followed a fairly high flux by an hour or so; and where the diurnal variation of  $f_oE$  on two days follows much the same course, except for the period 01-04 h, when an increase in  $f_oE$  on 12 November coincides with two high fluxes, while lower values on 10 November accompany lower fluxes.

Fig. 2 shows the  $f_oE$  values, with their scaling limits, for 03 h 00, together with the values of the directional electron flux recorded by Alouette I on the pass between 02 and 04 h. The similarity in shape of the graphs is striking and suggests that there is indeed a definite correlation between the two. Similar graphs for other times are equally encouraging and prompted the more thorough examination reported in the following sections.

## Correlation analysis

During November the solar zenith angle,  $\chi$ , at a given local time changes fairly rapidly as the month progresses, so that part at least of the general increase of  $f_oE$  towards

the end of the period in Fig. 2 is due to this cause. For the purpose of correlation, the values of  $f_oE$  at each hour of local time were therefore corrected to the value they would have had, if the solar zenith angle had been the same as it was in the middle of the period, on 9 November. These standard values of the solar zenith angle are given in the second column of Table 1.

To make this correction, the relationship between  $f_oE$  and  $\chi$  was first found, using only data for days and times when the flux of electrons measured by the satellite was low. When the solar zenith angle is near to 90°, as in our case, the relationship can be expressed in the form

$$f_oE = k [Ch(x,\chi)]^{-n} \quad (1)$$

where  $k$  and  $n$  are constants and  $Ch(x,\chi)$  is the "grazing incidence function" introduced by Chapman (1931) and tabulated by Wilkes (1954). Here  $x$  is the minimum radial distance from the centre of the earth to the ray from the sun to the point under consideration, expressed in atmospheric scale heights. At Argentine Islands the virtual heights of the E layer sometimes rise as high as 180 km during sunrise and sunset, and fall to about 110 km at midday. Since we are concerned with times when the solar zenith angle is close to 90°, we have chosen  $x = 150$  as a working value, corresponding to a scale height of 42 km. This agrees well with Fig. 15 of Nicolet (1960) but may be too large a scale height according to later work (Jacchia, 1971). The smallness of the correction which we are computing makes the approximation quite acceptable.

From a graph of  $\log(f_oE)$  against  $\log[Ch(x,\chi)]$  the

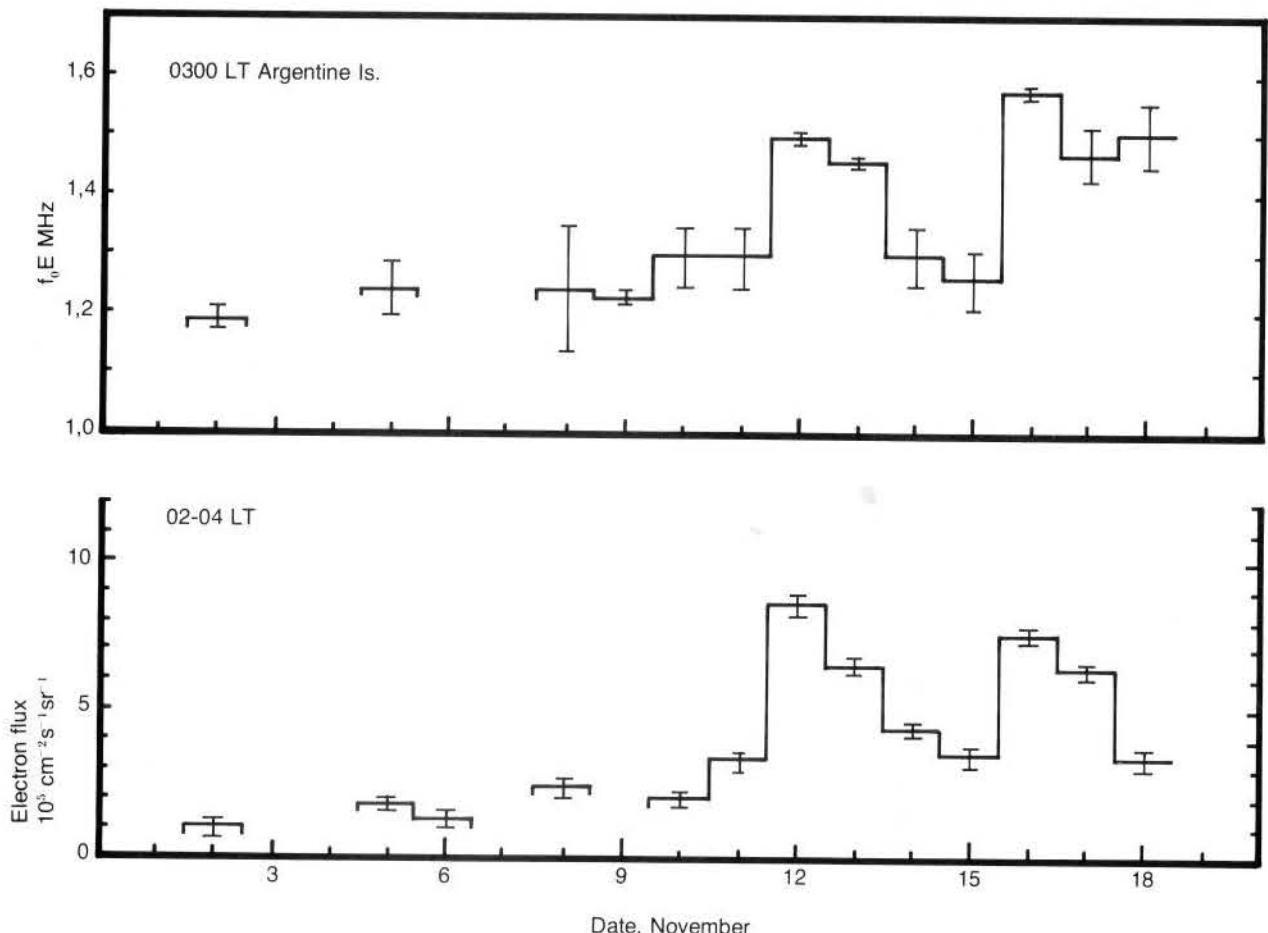


Fig. 2. Variation of  $f_oE$  at 03 h 00 at Argentine Islands and of electron flux observed by Alouette I between 02 h 00 and 04 h 00 in the vicinity.

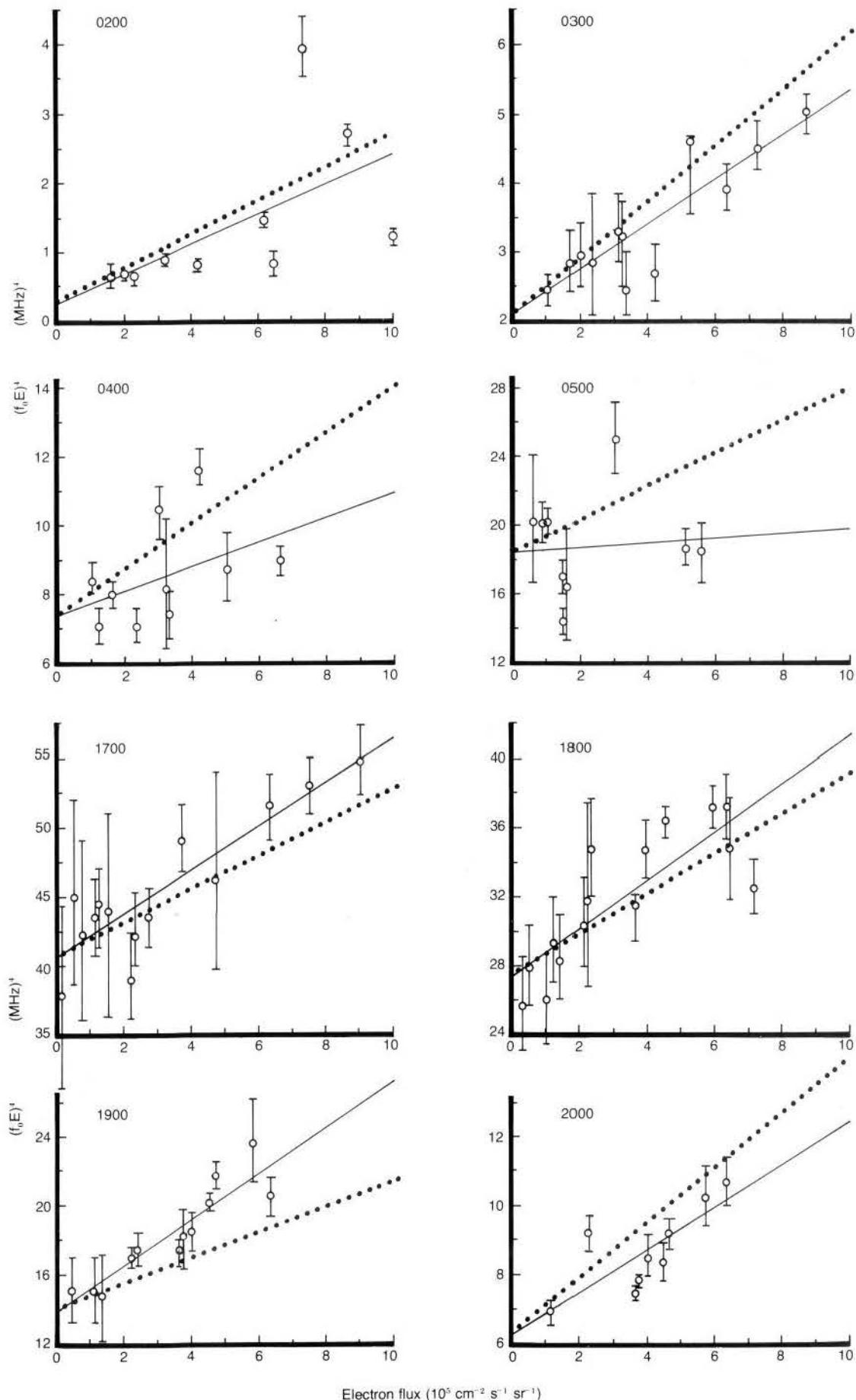


Fig. 3. Plots of  $(f_0 E)^4$ , corrected to a constant solar zenith angle at each local time, against the directional electron flux  $>40 \text{ keV}$  recorded within 1.5 hours by Alouette I. Full lines – least squares fitted straight lines; Dotted lines – theoretical straight lines calculated approximately from observed electron fluxes.

values of  $k$  and  $n$  in equation (1) were found to be 3,67 MHz and 0,36 respectively. Each observed value of  $f_0E$  was then corrected by subtracting from it the difference of the value calculated for that time and date and the value calculated for the standard solar zenith angle for the same time. The corrections were typically about 0,03 MHz and never exceeded 0,12 MHz. No corrections were applied to the data for 02 h 00, nor to those for 03 h 00 before 9 November, because  $\chi$  was greater than 90° under these conditions and  $Ch(x, \chi)$  then became so large and sensitive to changes in the scale height that it was thought better not to attempt the correction.

$f_0E$  has also been shown to vary with the Zürich sunspot number,  $R_z$ . Beynon & Brown (1959) have shown, for example, that there is on occasions a close correlation between day-to-day variations of  $f_0E$  and  $R_z$ . It seems doubtful, however, whether individual hourly values could properly be corrected to a standard sunspot number using a relationship derived for longer-period variations. Fortunately,  $R_z$  was very low, of the order of 10, for most of the period we are considering, and only approached 50 on three days. A calculation gave the increase of  $f_0E$  corresponding to  $R_z = 50$  as less than 0,05 MHz. It has therefore been assumed that the effect of sunspot number variations on our  $f_0E$  values is minimal and no correction has been made.

If the E region were ionized by solar electromagnetic radiation alone, the continuity equation for electrons would be written:

$$\frac{dN_o}{dt} = Q_o - \alpha N_o^2 \quad (2)$$

If  $dN_o/dt$  is negligible in comparison with  $Q_o$  and  $\alpha N_o^2$ , as it is in our case, we can write

$$Q_o = \alpha N_o^2 \quad (3)$$

Here  $Q_o$  is the production rate of electrons due to the sun,  $\alpha$  is the recombination coefficient and  $N_o$  is the number of electrons per  $\text{cm}^3$ , all at the E layer peak. If the precipitated electron flux produces  $Q_e$  ion pairs per  $\text{cm}^3$  per second, then, still assuming that  $dN/dt=0$ , the total electron density  $N$  is given by

$$Q = Q_o + Q_e = \alpha N^2 \quad (4)$$

Since  $N = 1,24 \times 10^4 (f_0E)^2$

we can combine equations (3) and (4) to get

$$(f_0E)^4 = (f_0E_0)^4 + Q_e/(1,54 \times 10^8 \alpha) \quad (5)$$

where  $f_0E_0$  is the critical frequency which the E layer would have in the absence of ionization by the precipitated electrons.

Now the production rate  $Q_e$  is proportional to the flux  $J$  of the precipitated electrons (Chamberlain, 1961). Thus if we plot  $(f_0E)^4$  against  $J$ , we should obtain a straight line, the intercept of which will be  $(f_0E_0)^4$ . The corrected values of  $(f_0E)^4$  are shown plotted against the average values of the electron flux during the nearest pass in Fig. 3. The error limits shown are derived from those estimated during the scaling of the ionograms. Ambiguous cases with several cusps of doubtful significance have been omitted.

There can be no doubt about the reality of the connection between the electron flux and  $f_0E$ . The correlation coefficient  $r$  is given in the third column of Table 1 and its significance level (David, 1938) in the fourth column.

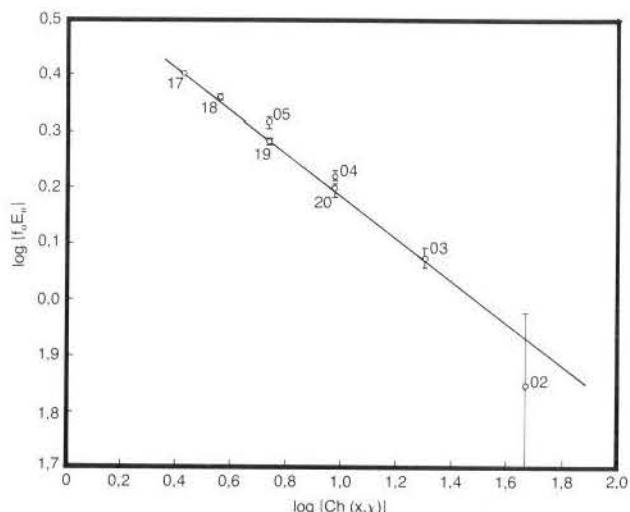


Fig. 4. Relationship between the logarithms of  $f_0E_0$  and the Chapman grazing incidence function  $Ch(x, \chi)$ .

The correlation is a little weak for 02 h 00, where the correction for solar zenith angle change could not be used, but breaks down for 04 h 00, where the significance level is only 0,2, and especially for 05 h 00, where it is completely insignificant. We are unable to give any explanation for this behaviour, which is surprising in view of the good correlation at the same solar zenith angles in the evening hours. It may be significant that there are, at both these times, only two data points corresponding to large fluxes, so that there may be insufficient data on which to base a conclusion.

The straight lines in Fig. 3 have been fitted by a least squares method and their slopes and intercepts are given in Table 1, together with the resulting values of  $f_0E_0$  and their estimated error limits, in this case derived from the standard deviations of  $(f_0E_0)^4$ . The logarithms of these values of  $f_0E_0$  are plotted against  $\log Ch(x, \chi)$  in Fig. 4. The straight line is a least squares fit, weighting the points inversely as the squares of the standard deviations in  $\log(f_0E_0)$ . It corresponds to  $k=3,614$  MHz and  $n=0,374$  in equation (1).

## Discussion

It would be interesting to compare the values of  $f_0E$ , as observed, and shown in Fig. 3, with the variations which we would expect theoretically from the measured electron fluxes. There are several reasons why this cannot be done with much confidence.

- (i) The electrons observed by Alouette I are in the energy range 40-250 keV, whereas the E region is most affected by electrons of energy about 4 keV.
- (ii) The instruments on Alouette I only observed the directional flux over a very limited range of pitch angles during one pass through the region of interest. Thus no information on the pitch angle distribution of the electrons at the time of the events is available.
- (iii) The only data available on the energy spectrum of the electrons are the fluxes  $> 40$  keV and  $> 250$  keV from the two counters on the satellite.

In the absence of this detailed information it is not possible to compare N-h profiles, calculated from the ionograms, with the profiles to be expected from the precipitated electron flux. This would clearly be the most satis-

factory way of proving the hypothesis of a connection between the two.

It was, however, possible to make a rough estimate of the effect in the following way. From the fluxes measured by the counter with a threshold of 40 keV and that with its threshold at 250 keV, corrected for background, a characteristic energy,  $E_o$ , was estimated, assuming that the spectrum could be represented by an equation of the type (Bailey, 1968)

$$J_{>E} = J_0 \exp(-E/E_o) \quad (6)$$

From an examination of the results for the period under investigation an average value of  $E_o = 50$  keV seemed appropriate. Further, Greener (1973) has examined the pitch angle distribution of the electrons observed by Alouette I, in the range of L values of interest here, and from his figures an average pitch angle distribution was found. With this information, the omnidirectional flux at 300 km could be estimated and then the method of Wulff & Gledhill (1974) was used to calculate the production rate,  $Q_e$ , at 120 km due to this flux. Assuming a value of  $2 \times 10^{-7} \text{ cm}^3 \text{ s}^{-1}$  for  $\alpha$  at this height (Bailey et al., 1970), the production rate  $Q_o$  was found from the values of  $f_o E_o$  in Table 1 and added to that due to the precipitated electrons to give the new value of  $(f_o E)^4$  using equation (4). The results of this rough calculation are shown by the dotted lines in Fig. 3. They are seen to agree far better than we have any reason to expect, considering the multitude of approximations in the calculation. They do, however, go some way towards confirming that the effects reported are in fact due to the electrons precipitated from the radiation belts.

Many criticisms may be levelled at details of the very simple approach which we have used in our analysis. For example, our use of the recombination coefficient  $\alpha$  as if the recombination were a simple square law process is surely incorrect even at 150 km. Again, the exponents n in equation (1) which we have evaluated, of the order of 0.37, show that the assumption that  $Q_o$  will be proportional to  $(f_o E_o)^4$  must be wrong. The increase of this value over the figure of 0.25 derived from simple theory is probably mainly due to the temperature gradient in the vicinity of 150 km (Rishbeth & Garriott, 1969). Nevertheless, the changes in  $f_o E$  with which we are dealing are in fact so small that it makes little difference what power is plotted against the electron flux; simple graphs of  $f_o E$  against flux give good straight lines, and extrapolated values for  $f_o E_o$  which are almost identical with those found in Table 1. Therefore we believe that our work does demonstrate the effect of the precipitated electrons in a way which merits confidence.

## Conclusions

The increases of  $f_o E$  reported in this paper are relatively small, of the order of 0.3 MHz for an observed flux  $>40$  keV of  $10^6$  electrons  $\text{cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$ , but, as the error bars show, they can be accepted with some confidence as real variations. As equations (4) and (5) show, even a production rate  $Q_e$ , due to precipitated electrons, equal to that due to solar radiation,  $Q_o$ , would only increase  $f_o E$  by 20%. In other words,  $f_o E$  is not a sensitive indicator of changes in the ion production rate. Looked at from this point of view, the ionization rate due to precipitated electrons exceeds the solar rate by factors which may be

greater than 100 at 02 h 00, often exceeds the solar rate at 03 h 00, and may equal it when the flux is high even at 04 h 00, 19 h 00 and 20 h 00. The importance of electrons as ionizing agents in the E region is certainly not negligible over these periods at Argentine Islands. Similar analyses at other stations would show whether this is a consequence of the station's proximity to the South Atlantic Anomaly or whether the phenomenon is more widespread, as the work of Ivanov-Kholodny & Kazatchevskaya (1971), for example, suggests.

Clearly the method used here is capable of very considerable elaboration. With the availability of modern soft electron spectra and pitch angle information, such as that provided by the LEPEDEA detectors flown on several satellites by Frank (1967) and the instruments used by Heikkila et al. (1970) on the ISIS satellites, it should be possible to make detailed comparisons of the N-h profiles derived from ionograms, both topside and ground-based, with the effects to be expected on the basis of the energy and pitch angle spectra of the precipitated electrons as simultaneously observed nearly overhead. This aspect of the work is currently being pursued.

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